

Effect of Defect Density Variation in CdTe Solar Cell Using Cu as A Hole Transport Layer and Back Contact

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Abstract – Cadmium Telluride (CdTe) solar cells have the potential to achieve high efficiency, but material defects present a significant challenge to their performance. These defects act as recombination sites, capturing electrons and holes, thereby hindering current flow and reducing the overall efficiency of the CdTe solar cells. Addressing these defects is crucial for enhancing charge carrier transport and optimizing efficiency. This research paper analyzes the effect of defect density numerically using SCAPS-1D simulation software on SnO₂/CdTe/CdS/Cu multi-junction structural solar cells. The performance parameters of the proposed cell are optimized by varying different numerical particulars. Our analysis reveals that by varying the defect density from 10¹² to 10¹⁸ cm⁻³, the efficiency of the proposed solar cells varies from 17.65% to 2.56%, with the open-circuit voltage (Voc) at 0.85V, short-circuit current density (Jsc) at 26 mA/cm², and fill factor (FF) at 82%. The efficiency optimization is attributed to the conduction band effective density of states, Nc = 2.2 × 10¹⁸ cm⁻³, and valence band effective density of states, Nv = 1.8 × 10¹⁹ cm⁻³ of the donor material (SnO₂). The cell's power conversion efficiency is optimized when the absorber has a balanced concentration of neutral, acceptor, and donor defects at 1 × 10¹⁴ cm⁻³. At an operating temperature of 300 K, an efficiency of 18% is achieved.

Keywords: Cadmium telluride (CdTe) solar cell, Defect density, Recombination sites, Solar cell performance

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I. Introduction

Research on CdTe solar cells has been ongoing since the mid-20th century to improve their efficiency for enhanced solar energy utilization [1]. These cells are of great interest due to their excellent performance. Recent studies have shown that their efficiency can exceed 22% [2] – [4]. CdTe solar cells are considered to be one of the most promising solar cell technologies for the future. They are relatively inexpensive to manufacture due to abundant raw materials, simple manufacturing processes, thin-film technology requiring less material, and scalable production and have a high absorption coefficient of 5 × 10⁵ cm⁻¹, which means that they can absorb a large amount of sunlight approximately 400 to 800 nanometers (nm). The direct band gap of CdTe (1.5 eV) is also well-matched to the solar spectrum, which means that it can absorb photons from a wide range of wavelengths. This makes CdTe solar cells a very efficient way to convert sunlight into electricity [5] – [6]. Incorporating desirable attributes of CdTe's absorption characteristics, the presence of n-type material is imperative for the optimal formation of a

high-quality p-n junction. Notably, Cadmium Sulphide (CdS) stands out as a prominent candidate, showcasing exceptional capabilities due to its direct band-gap of approximately 2.42 eV within the visible spectrum.

The transparent nature of CdS enables solar irradiation to transmit through the CdTe layer, thereby enhancing the photovoltaic response. [7] – [8]. Conversely, the heterojunction CdTe/CdS solar cell necessitates the integration of a transparent conducting oxide (TCO) layer at its front contact. This strategic inclusion allows photons to enter the cell, with the active interface acting as an electrode to efficiently collect the generated current. [9] – [11]. Notably, transparent conducting oxide (TCO) layers commonly employed for this purpose include Fluorine-doped Tin Oxide (SnO₂:F or FTO) and Aluminum-doped Zinc Oxide (ZnO: Al or AZO) [12] – [13]. These materials play a pivotal role in enabling photon transmission and effectively harnessing current at the active interface of the heterojunction CdTe/CdS solar cell [14] – [15]. Incorporating a high-resistive transparent (HRT) film

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between the transparent conducting oxide (TCO) and the CdS buffer layer holds the potential to enhance efficiency by mitigating the impact of non-uniformity [16]. This approach effectively minimizes variations that may arise, contributing to a more consistent and optimized performance of the solar cell [17]. Undoped Tin Oxide ($i\text{-SnO}_2$) and undoped Zinc Oxide ($i\text{-ZnO}$) stand out as materials frequently employed for the development of a high-resistive transparent (HRT) layer [18, 19]. This layer serves a crucial role in enhancing the solar cell's performance by reducing the impact of non-uniformity, thereby contributing to a more consistent and efficient operation [20] – [21].

During the 1980s, T. P. Shalvey et. al. and others focused on the back surface field (BSF) layer and its consequential influence on the efficiency of solar cells [22] – [24]. These studies delved into the effects and potential benefits of implementing a BSF layer in photovoltaic devices, aiming to optimize energy conversion and enhance overall solar cell performance [25] – [26]. A material with a substantial band gap is strategically used as a barrier at the CdTe/BSF heterojunction [27] – [28]. This barrier effectively prevents carriers from dissipating at the back contact interface [29]. By minimizing carrier losses, this approach enhances the overall efficiency of the photovoltaic system [30].

Furthermore, this approach mitigates carrier recombination at the back contact interface, resulting in an enhancement of cell efficiency [31] – [32]. By curbing the tendency for carriers to recombine at this junction, the overall performance of the solar cell is notably improved [33] – [34]. Within the same scholarly framework, investigations have established that the utilization of a hole transport-electron blocking layer (HT-EBL) parallels the role fulfilled by the back surface field (BSF) layer [35] – [36]. This layer, the HT-EBL, effectively reduces carrier losses and mitigates recombination phenomena at the rear contact interface [37]. As a result, the discernible consequence manifests as an enhancement in the efficiency of solar cells. Additionally, it substantiates a significant contribution toward the elevation of the solar cell's operational efficacy. A notably promising candidate material employed as a hole transport-electron blocking layer (HT-EBL) within CdTe solar cells is cuprous oxide (Cu_2O) [38, 39]. A notably promising material used as a hole transport-electron blocking layer (HT-EBL) in CdTe solar cells is cuprous oxide (Cu_2O) [37] – [38]. With a wide band gap ranging from 2.1 eV to 2.61 eV and natural p-type conductivity, cuprous oxide has characteristics that make it suitable for this purpose [39] – [40].

Cu_2O is non-toxic, making it environmentally friendly. It is also readily available, cost-effective, and has a high absorption coefficient in the visible spectrum, making it a good choice for solar cells [40]. However, defects in CdTe solar cells can limit their efficiency. Common defects

include grain boundaries, impurities, and dislocations, which can create recombination centers for charge carriers, reducing the overall performance of the cell [41].

In the proposed $\text{SnO}_2/\text{CdS}/\text{CdTe}/\text{Cu}$ model, copper (Cu) serves a dual purpose as both the back contact and an electrode, thereby facilitating the accumulation and directed flow of generated electrons from the CdTe layer into the external circuit. This orchestrated movement of electrons forms the fundamental basis for the overarching process of electricity generation. It is pertinent to underscore that the inclusion of copper at the back contact interface exerts a notable influence on carrier recombination phenomena. Acting as an effective electron-blocking layer, the Cu layer prevents photo-generated electrons from recombining at the rear surface. This process increases the retention of generated charge carriers, thereby improving the efficiency of energy conversion.

In summary, the strategic addition of copper in the $\text{SnO}_2/\text{CdS}/\text{CdTe}/\text{Cu}$ model has two main effects: it enhances charge collection and reduces carrier recombination. This combination leads to an overall increase in the efficiency of the photovoltaic device.

II. Solar Cell Structure and Materials Properties

Fig.1 depicts the schematic configuration of a $\text{SnO}_2/\text{CdS}/\text{CdTe}/\text{Cu}$ heterojunction solar cell. The cell primarily consists of a $\text{CdS}(n)/\text{CdTe}(p)$ heterojunction, with n-type and p-type doping concentrations of $1 \times 10^{18} \text{ cm}^{-3}$ and $2 \times 10^{14} \text{ cm}^{-3}$, respectively, forming the core structure. An optically transparent SnO_2 film functions as the high-resistive transparent (HRT) layer situated between the transparent conducting oxide (TCO) and CdS regions. The back contact is established through a copper (Cu) layer. Key constituents include CdS and CdTe layers, serving as the window and absorber materials, respectively. Cu and FTO materials fulfill distinct roles as the rear contact (anode) and the front contact (cathode), respectively. The FTO layer on the glass substrate is characterized by a thickness of 1 mm, a refractive index of 1.52, and a front surface reflection coefficient of 4.22%. In this context, parameters such as thickness, carrier concentration, interface defect density, and operational temperature have been meticulously chosen as focal considerations.

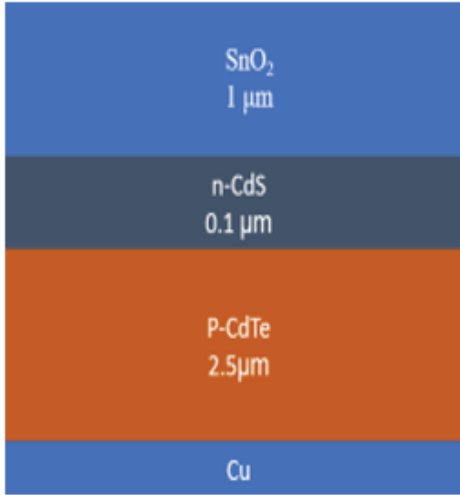


Fig. 1. FTO/CdS/CdTe/Cu solar cell model

To facilitate simulation, the glass/TCO substrate is modeled as an FTO layer, specified with a thickness of 300 nm [42] and an n-type doping concentration of $3 \times 10^{20} \text{ cm}^{-3}$. Comprehensive device geometrical and physical attributes are shown in Table I for reference and analysis. The absorber layer in solar cells is made up of CdTe, which has a high absorption coefficient of over 10^4 cm^{-1} and an ideal direct band gap of 1.45 eV for effective photo conversion at room temperature [43] – [44]. To form a p-n junction, p-type CdTe is often paired with n-type CdS, which serves as the window layer in the heterojunction [45] – [46]. The absorber layer may contain acceptor defects because of its p-type characteristics, which can lead to a reduction in solar cell efficiency. CdS serves as the window layer and is commonly used in CdTe photovoltaic devices [46]. This layer facilitates Cadet Nucleation and growth and creates a high-quality heterojunction with the absorber shown this band alignment from the energy band diagram on SCAPS [47]. At room temperature, CdS has a direct bandgap of 2.42 eV and exhibits an acceptable lattice mismatch of 9.7% with CdTe [48]. This layer is n-type so donor defects are found in this layer. The band gap of CdS, measuring 2.42 eV, signifies its capability to facilitate relatively unimpeded transmission of a significant portion of the solar spectrum. This attribute renders it an apt choice as a window layer situated atop the absorber. Herein, the occurrence of photo-generated surplus minority carriers finds an advantageous environment for its manifestation, a concept underscored by Coutts (1982). In the context of CdTe, it manifests as a direct band gap semiconductor at 1.5 eV. To construct solar cells based on CdTe, films of this material measuring 2.5 μm in thickness are employed [47].

TABLE I
DEVICE PARAMETERS USED IN THE SIMULATION

| Parameters | Layers | | | |
|--|----------------------|----------------------|----------------------|----------------------|
| | CdTe | CdS | SnO ₂ | FTO [17] |
| Thickness, t_m (nm) | 2500 [47] | 100 | 1000 | 300 |
| Electron affinity, χ_e (eV) | 3.9 | 10 | 4 | 3.2 |
| Dielectric ratio, ϵ_r | 9.4 | 4 | 9 | 4.4 |
| Band gap, E_g (eV) | 1.45 | 2.42 | 3.63 | 9 |
| Conduction band effective density of states, N_c (cm^{-3}) [17] | 8×10^{17} | 2.2×10^{18} | 2.2×10^{18} | 2.2×10^{18} |
| Valance band effective density of states, N_v (cm^{-3}) | 1.8×10^{19} | 1.8×10^{19} | 1.8×10^{19} | 1.8×10^{19} |
| Electron mobility, μ_n (cm^2/Vs) | 320 | 100 | 100 | 90 |
| Hole mobility, μ_p (cm^2/Vs) | 40 | 25 | 25 | 90 |
| Carrier concentration, N_a (cm^{-3}) | 2×10^{14} | 0 | 0 | 3×10^{20} |
| Carrier concentration, N_d (cm^{-3}) | 0 | 1.1×10^{18} | 1×10^{17} | 0 |
| Hole lifetime (ns) | 1 | 0.01 | 0.1 | 1×10^{15} |
| Electron lifetime (ns) | 1 | 0.01 | 100 | |

The collective assortment of parameters crucial for the design of the SnO₂/CdS/CdTe/Cu configuration is comprehensively shown in Table I. This authors have chosen a 0.25 cm^2 area for this device. The simulation was conducted with an incident irradiation of 1000 W/m^2 .

The investigation at hand employs SCAPS-1D for modeling and numerical simulation of the

SnO₂/CdS/CdTe/Cu heterojunction solar cell. SCAPS-1D effectively facilitates the analysis of the J–V (current-voltage) characteristics and quantum efficiency of the solar cell. This software affords a comprehensive platform for discerning and quantifying the performance aspects of the studied system. The utilization of SCAPS-1D involves a meticulous exploration and comparison of defects within CdTe solar cells, featuring similar parameters albeit distinct layers (CdTe & CdS). Given the crystalline nature of CdS and CdTe, their physical attributes exhibit variability contingent upon growth conditions and fabrication techniques employed. Within the scope of this study, it is pivotal to underscore that the parameters of the absorber and window layers remain constant, aligning precisely with the values outlined in Table I. This consistency is upheld throughout all simulation steps, ensuring robustness and accuracy. The parameters requisite for all four materials under examination have been diligently furnished. Meanwhile, a copper layer boasting a work function of 5 eV serves as the back contact, an operating temperature of 300K, and a total number of simulations have been carried out with a series resistor of 3.2 Ω and a shunt resistor of 3200 Ω.

III. Results and Discussion

We conducted an analysis focusing on two distinct types of defects: acceptor and donor—present within the CdTe and CdS layers. Our assessment encompassed the calculation of various critical parameters, encompassing efficiency, open-circuit voltage (Voc), short-circuit current density (Jsc), and fill factor (FF), employing the SCAPS platform. The range of defect densities scrutinized extended from 10¹² to 10¹⁸. It emerges that acceptor defects are predominant within the CdTe layer, whereas donor defects are primarily concentrated within the CdS layer. Notably, defect densities beneath 10¹² cm⁻³, exhibited minimal impact on the percentage of the power conversion efficiency (PCE) for both layers. However, a discernible reduction in all parameters was observed as defect density exceeded the 10¹² cm⁻³, threshold. This

decrement was found to be more pronounced in the CdTe layer compared to the CdS layer. At a defect density of 10¹⁸ cm⁻³, the lowest values for Voc, Jsc, FF, and PCE were identified. These outcomes collectively underscore the correlation between heightened defect density and diminished solar cell performance. The findings validate the notion that the augmentation of defect density invariably precipitates a decline in the operational efficiency of the solar cell.

A. Effect of the Acceptor Defect Density within the CdTe Absorber Layer

Our simulation outcomes indicate that Acceptor defects have a greater impact on cell performance. In Fig. 2, it is clear that acceptor defects elicit a profound transformation across all parameters. Efficiency precipitously declines from 17.65% to 2.56%, and the fill factor undergoes a substantial reduction, descending from 84.77% to 49.87%. This discernible divergence is rooted in the energy band diagram for the CdTe solar cell incorporating defects. Our analysis vividly elucidates that the introduction of defects engenders trap levels. The nature of these trap levels evolves, transitioning from lying beneath the Fermi level of electrons for lower defect densities to situating between the Fermi level of electrons and the conduction band energy level at a defect density of 10¹⁸ cm⁻³. Consequently, the phenomenon of photo-generated electrons becoming ensnared at these energy levels ensues, thereby manifesting as a diminishment in the photovoltaic efficacy of CdTe solar cell devices.

Another explanatory facet resides in the absorber layer's role. The augmentation of efficiency correlates closely with the characteristics of the absorber layer. Consequently, the influence of increased defect density is more pronounced within the absorber layer, further substantiating the observed dissimilarity in the impact of acceptor and donor defects.

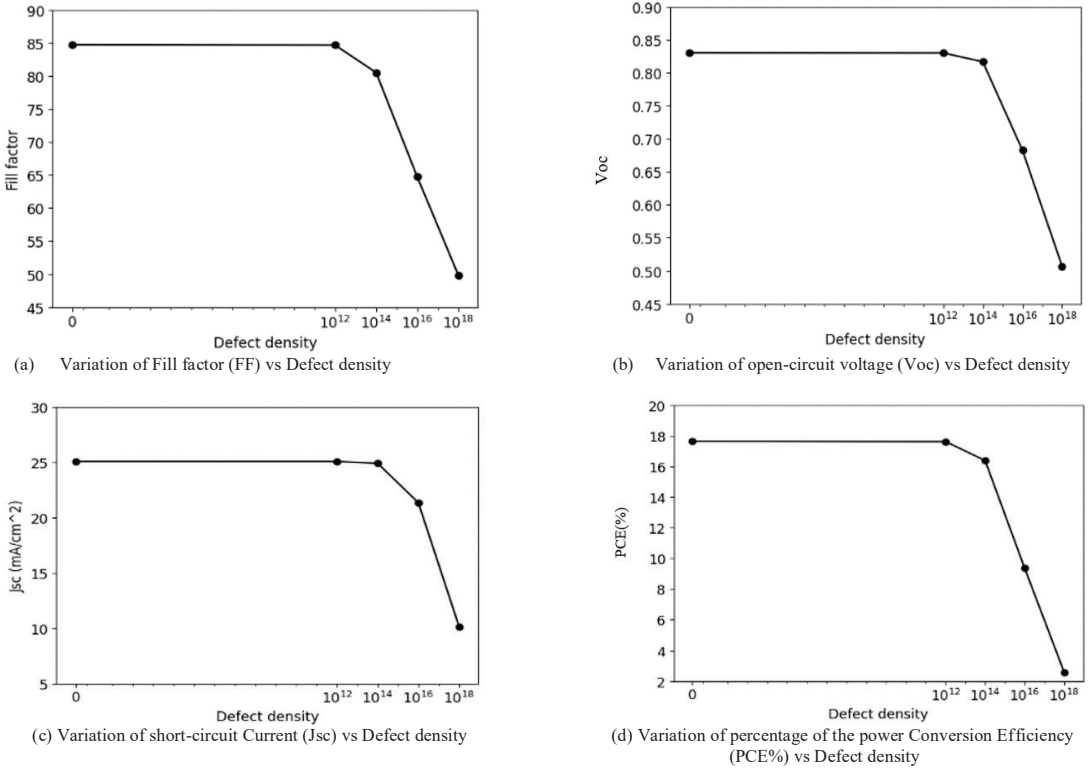


Fig. 2. Effect of the acceptor defect density on FF, J_{sc} , V_{oc} , and PCE(%) in the CdTe absorber layer.

B. Effect of the Donor Defect Density in CdS Window Layer

In our observations, a notable trend emerges regarding the impact of defect density on efficiency within the CdTe absorber layer. The efficiency demonstrates a marked reduction, declining from 17.65% to 2.56% as the defect density escalates from 10^{12} to 10^{18} cm^{-3} , as shown in Fig. 2. This phenomenon is distinct in the context of acceptor defects.

Similarly, in the case of the CdS Window layer, a distinct relationship between defect density and efficiency is established. Here, the efficiency exhibits a discernible

decline from 16.56% to 13.82% as the defect density increases from 10^{12} to 10^{18} cm^{-3} , as shown in Fig. 3. This correlation underscores the influence of defect density on the performance of the solar cell in the context of both layers. The effect of different densities and defects charge types in the light-absorbing CdTe & CdS layer on the photovoltaic performance parameters of solar cell devices has been considered. It has been found that photovoltaic parameters retain their highest values of V_{oc} , J_{sc} , FF, and η till a density of defects of 10^{12} cm^{-3} . This research indicates that there is an optimum defect density of 10^{12} cm^{-3} , above which the solar cell performance starts degrading.

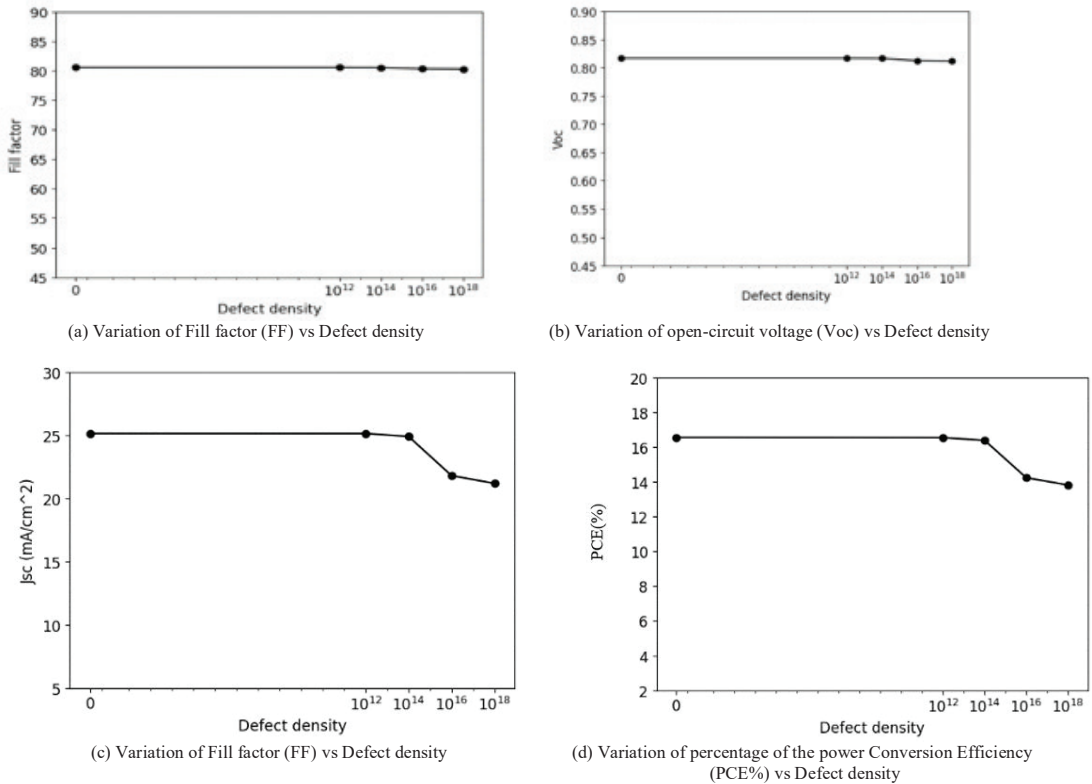


Fig. 1. Effect of the donor defect density on FF, Jsc, Voc, and PCE (%) in the CdS window layer.

C. Spectral response

The concept of Quantum Efficiency (QE) revolves around the proportion of incident photons that are effectively captured by a solar cell's total charge carriers. In this investigation, we have generated four distinct QE curves shown in Fig. 4, with curves (a) and (c) delineating acceptor defects, while curves (b) and (d) portraying donor defects. These curves vividly illustrate the defect characteristics inherent to both the absorber and window layers. Examining the QE curve attributed to acceptor defects within the CdTe layer, a noticeable trend emerges wherein efficiency experiences a rapid decline in tandem with escalating defect density. This phenomenon is eloquently portrayed in Fig. 4. (a), which showcases the generated curves labeled 1(black), 2(red), 3(blue), and 4(green). Notably, curves 1 and 2, corresponding to defect densities of 10^{12} and 10^{14} cm⁻³, exhibit a relatively similar efficiency drop. In contrast, curves 3 and 4, corresponding to defect densities of 10^{16} and 10^{18} cm⁻³, demonstrate a substantial reduction in efficiency. An intriguing nuance surfaces beyond a wavelength of 500, where curve 4

notably flattens. This is attributed to heightened recombination and a concomitant decline in carrier collection, culminating in a noticeable reduction in quantum efficiency. In essence, these QE curves proficiently unveil the intricate interplay between defect density, carrier behavior, and quantum efficiency, shedding light on the dynamic response of the CdTe layer to varying fault densities.

The QE curve corresponding to donor defects is portrayed in Fig. 4. (b). This graph delineates a decline in efficiency as defects escalate, albeit at a slower pace compared to the absorber layer. This specific depiction pertains to the CdS layer's QE curve. Notably, the produced curves labeled 1, 2, 3, and 4 correspond to defect densities of 10^{12} , 10^{14} , 10^{16} , and 10^{18} cm⁻³, respectively. It is intriguing to note that, akin to the acceptor defect case, curves 1 and 2 bear a similarity due to a relatively marginal efficiency drop between defect densities 10^{12} and 10^{14} cm⁻³. However, noteworthy distinctions emerge when comparing curves 3 and 4 to curves 1 and 2. Notably, the size of donor defects being relatively small, the efficiency decrement is less pronounced as defect density increases. Moreover, the influence of recombination initially

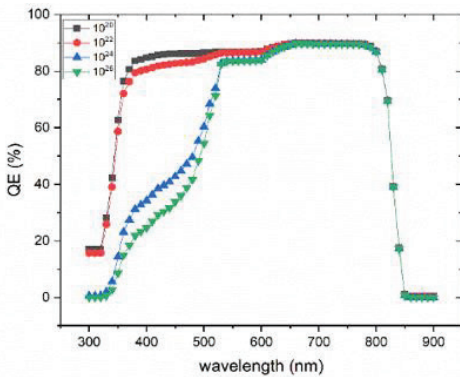
intensifies for curves 3 and 4, followed by a subsequent decline in the later wavelengths. Consequently, curves 3 and 4 exhibit a characteristic dip and recovery in efficiency.

The delineation of the device's external quantum efficiency (QE) effectively captures the solar cell's functional wavelength range, thereby enhancing our insight into the operational dynamics of the investigated solar cell. This figure of merit is defined as the ratio between the effective photocurrent calculated through the cell and the source photocurrent which is in the form of [48] as in (1),

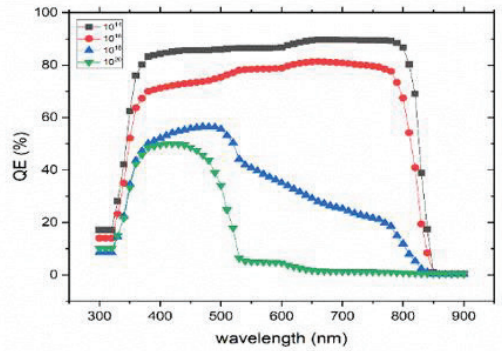
$$I_{src} = \frac{1}{hc} q \lambda P_b W_b \tag{1}$$

where λ is the photon wavelength, q is the electronic charge, c is the speed of light, h is Planck's constant P_b is the power density of the incident light, and W_b is the beam width clipped to the device [49].

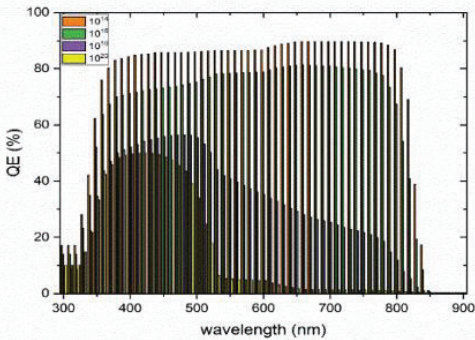
As dictated by a given light source, the available photocurrent across the cell is intrinsically contingent upon the distinct optical attributes of the various materials at play. Mainly, we have adopted the wavelength-dependent refractive index (n) and extinction coefficient (k) of both CdS and CdTe layers, as elucidated by experimental extraction within the 250–1500 nm wavelength spectrum, as detailed in. Similarly, the optical characterization of the SnO₂ region was drawn.



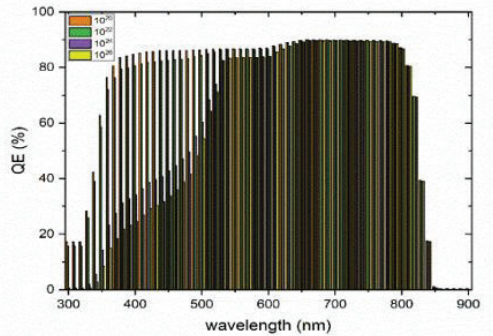
a) Variation of QE (%) vs wavelength for acceptor defect in CdTe Layer.



b) Variation of QE (%) vs wavelength for Donor defect in CdS Layer.



c) Variation of QE (%) vs wavelength for acceptor defect in CdTe Layer by Bar curve



d) Variation of QE (%) vs wavelength for Donor defect in CdS Layer By Bar curve

Fig. 2. Dependency of QE on defect density.

In Fig. 4, the external quantum efficiency (QE) profiles of both the reference and optimized cells are shown in detail. The QE behavior is demonstrated across a wavelength range from 300 nm to 950 nm. It is notable that the optimized configuration shows a clear increase in QE across the entire explored wavelength spectrum. This trend continues until the CdTe optical bandgap causes an absorption edge, which occurs around 850 nm. Particularly noteworthy is the peak of absorption within the range of $500 < \lambda < 800$ nm, where the average QE increases to approximately 96%, reaching a peak of 96.98% precisely at $\lambda = 650$ nm. The distinct absorption point at $\lambda = 500$ nm is a result of the inherent properties of the CdS material. It is also worth noting that for wavelengths below this critical value, light is primarily absorbed by the FTO layer. This comprehensive illustration of the QE behavior and its interaction with the optimized cell design is now complete.

IV. Conclusion

This study offers a thorough analysis of how defects impact the performance of CdTe solar cells. The focus is on acceptor and donor defects in the CdTe absorber and CdS window layers, respectively. Using the SCAPS platform, we calculated key parameters like efficiency, open-circuit voltage (Voc), short-circuit current density (Jsc), and fill factor (FF) across a range of defect densities from 10^{12} to 10^{18} cm⁻³. Our findings show that acceptor defects are more prevalent in the CdTe layer, while donor defects primarily affect the CdS layer. Our simulations demonstrate that defect densities below 10^{12} cm⁻³ have minimal impact on the efficiency of both layers. However, as defect density increases beyond this threshold, all parameters, including Voc, Jsc, FF, and overall efficiency, exhibit a significant decline. The CdTe layer is more affected, where efficiency drops from 17.65% to 2.56% as acceptor defect density increases from 10^{12} to 10^{18} cm⁻³. In contrast, donor defects in the CdS layer result in a smaller efficiency reduction, from 16.56% to 13.82%. The introduction of acceptor defects in the CdTe layer leads to the formation of trap levels, which evolve with increasing defect density. These traps, initially below the Fermi level, eventually lie between the Fermi level and the conduction band at higher densities, causing photo-generated electrons to become trapped and reducing photovoltaic efficiency. The absorber layer's characteristics are crucial, as efficiency is closely tied to its quality; hence, defects in this layer have a more pronounced impact. Spectral response analysis through Quantum Efficiency (QE) curves further elucidates the effects of defect density. The QE for acceptor defects in the CdTe layer shows a rapid efficiency decline with increasing defect density, especially beyond a wavelength of 500 nm, due to increased recombination and reduced carrier collection. For donor defects in the CdS layer, the decline is slower and less pronounced, with a characteristic dip and

recovery in efficiency. This study highlights the critical role of minimizing defect densities in CdTe solar cells to optimize their performance. It emphasizes the need for improved materials and manufacturing processes to mitigate defects, particularly in the CdTe absorber layer, to enhance the efficiency and reliability of solar cell devices. Overall, this work contributes to a deeper understanding of the interplay between defect density and solar cell performance, providing valuable insights for future research and development in photovoltaic technology.

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Conflict of Interest

The authors declare no conflict of interest in the publication process of the research article.

Author Contributions

Author 1: Data collection, analysis, writing – original draft preparation; Author 2: Supervision, draft review and editing, investigation; Author 3: Conceptualization, review; Author 4: Project administration; Author 5: Analysis, writing – original draft preparation.

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